

To solve the wave eqn in polar... we will need BESSELS FUNCTIONS.

BESSELS FUNCTIONS - are solutions to Bessels Differential Equation.

$$x^2 y'' + x y' + (x^2 - p^2) y = 0$$

where p is a constant.

This ODE arises often when solving problems of $\nabla^2 u$ in polar coordinates.

To do this we use a special technique in odes

METHOD of FROBENIUS

OUTLINE SOLV: $A(x)y'' + B(x)y' + C(x)y = 0$

* NOTE... This is a special case of the POWER SERIES METHOD.

by assuming $y = \sum_{m=0}^{\infty} c_m x^{m+r}$

For this to work we need to assume some things about $A(x)$, $B(x)$, $C(x)$.

When $A(x) \neq 0$ for any x in the domain
 $y = \sum_{n=1}^{\infty} c_n x^n$
plug in find c_n .

$$\lim_{x \rightarrow x_0} (x - x_0) \frac{B(x)}{A(x)}$$

$$\lim_{x \rightarrow x_0} (x - x_0)^2 \frac{C(x)}{A(x)}$$

both finite.

Bessel's Eqn $A(x) = x^2$
 $B(x) = x$
 $C(x) = (x^2 - p^2)$

problem at $x=0$
 $A(x)=0$ and we "look" over y'' .

method of Frobenius will work if...

$$\lim_{x \rightarrow 0} x \frac{x}{x^2} = \lim_{x \rightarrow 0} 1 = 1 \text{ Finite } \checkmark$$

$$\lim_{x \rightarrow 0} (x)^2 \frac{(x^2 - p^2)}{x^2} = \lim_{x \rightarrow 0} x^2 - p^2 = -p^2 \text{ Finite } \checkmark$$

Plug $y = \sum_{m=0}^{\infty} C_m x^{m+r}$ into Bessel's ODE

$$y' = \sum_{m=0}^{\infty} C_m (m+r) x^{m+r-1}$$

$$y'' = \sum_{m=0}^{\infty} C_m (m+r)(m+r-1) x^{m+r-2}$$

- we assume here that y converges on the open interval I so that y is differentiable on I and the power series can be directly differentiated.

so our ODE becomes:

$$\sum_{m=0}^{\infty} C_m (m+r)(m+r-1) x^{m+r} + \sum_{m=0}^{\infty} C_m (m+r) x^{m+r} +$$

$$\sum_{m=0}^{\infty} C_m x^{m+r+2} - p^2 \sum_{m=0}^{\infty} C_m x^{m+r} = 0$$

shift the index $= \sum_{m=2}^{\infty} C_{m-2} x^{m+r}$

same form

this helps us gather terms!
 $x^r, x^{r+1}, x^{r+2}, \text{etc.}$

$$\sum_{m=0}^{\infty} C_m [(m+r)(m+r-1) + (m+r) - p^2] x^{m+r} + \sum_{m=2}^{\infty} C_{m-2} x^{m+r} = 0$$

now x^r terms are $m=0$, x^{r+1} are $m=1$, etc and the coefficients of each of these must go to zero.

$$m=0 \quad C_0 [r(r-1) + r - p^2] x^r = C_0 [r^2 - p^2] x^r = 0 \quad \text{so } C_0 = 0 \text{ or } r^2 - p^2 = 0$$

$$m=1 \quad C_1 [(1+r)(r) + (1+r) - p^2] x^{1+r} = C_1 [r^2 + 2r + 1 - p^2] = 0$$

$$\text{so } C_1 = 0 \text{ or } r^2 + 2r + 1 - p^2 = 0$$

$$m \geq 2 \quad C_m [(m+r)(m+r-1) + (m+r) - p^2] + C_{m-2} = 0$$

$$C_m [(m+r)^2 - p^2] + C_{m-2} = 0 \quad \text{or}$$

$$C_m = \frac{-C_{m-2}}{[(m+r)^2 - p^2]}$$

OK First... we don't want our leading term in the series to go to zero!

$C_0 \neq 0$ so this relates r and p by the characteristic equation $r^2 - p^2 = 0$
or $(r+p)(r-p) = 0$

$$\text{or } r = \pm p$$

Second if $(r^2 - p^2) = 0$ then

$C_1 = 0$ can't avoid this.

The rest are fixed through the recurrence relation

$$C_m = -\frac{C_{m-2}}{[(m+r)^2 - p^2]}$$

What do we have so far... two linearly independent solns of the form

$$y = \sum_{m=0}^{\infty} C_m x^{m+r}$$

for $r = \pm p$ and $C_1 = 0$

$$C_m = -\frac{C_{m-2}}{[(m+r)^2 - p^2]}$$

note C_0 still undefined
and C_m for $m = \text{even}$ only.

CASE: $r = p$
then

$$C_0 = C_0$$

$$C_m = -\frac{C_{m-2}}{[m^2 + 2pm]} = -\frac{C_{m-2}}{m[m+2p]}$$

$$C_2 = \frac{-C_0}{2(2+2p)} = \frac{-C_0}{2^2(1+p)}$$

$$C_4 = \frac{-C_2}{4(4+2p)} = \frac{-C_2}{2^2 \cdot 2(2+p)} = \frac{C_0}{2^4 \cdot 2(2+p)(1+p)}$$

$$C_6 = \frac{-C_4}{6(6+2p)} = \frac{-C_4}{2 \cdot 6(3+p)} = \frac{C_0}{2^6 \cdot 2 \cdot 3(3+p)(2+p)(1+p)}$$

look for the pattern:

$$C_{2m} = \frac{(-1)^m C_0}{2^{2m} m! (p+1)(p+2)\dots(p+m)}$$

so $y_1(x) = C_0 \sum_{m=0}^{\infty} \frac{(-1)^m x^{2m+p}}{2^{2m} m! (p+1)(p+2)\dots(p+m)}$ This is one linearly independent soln.

CASE: $r = -p$ then

$$y_2(x) = b_0 \sum_{m=0}^{\infty} \frac{(-1)^m x^{2m-p}}{2^{2m} m! (-p+1)(-p+2)\dots(-p+m)}$$

these converge for all $x > 0$ and these are Bessel's Functions.

We often rewrite these using the gamma function.
and give them their own variable:

$$J_p(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(p+m+1)} \left(\frac{x}{2}\right)^{2m+p}$$

BESSEL'S FUNCTIONS
OF THE FIRST KIND

$$J_{-p}(x) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(-p+m+1)} \left(\frac{x}{2}\right)^{2m-p}$$

so our soln should be written

$$y = C_1 J_p(x) + C_2 J_{-p}(x)$$

Now... these solutions are only linearly independent when $p \neq \text{integer}$.

Ex. $J_{1/2}(x) = \sqrt{\frac{2}{\pi x}} \sin(x)$ $J_{-1/2}(x) = \sqrt{\frac{2}{\pi x}} \cos(x)$

BUT if p is an integer then $J_{-p}(x) = (-1)^p J_p(x)$
and we actually need to do more work
to find a linearly independent soln!

If we did this it would be complicated but we could show that.

$$Y_p(x) = \frac{J_p(x)\cos(\pi p) - J_{-p}(x)}{\sin(\pi p)}$$

BESSELS FUNCTIONS
OF THE SECOND KIND

GRAPHS of J_p and J_{-p} and Y_p on computer...
in files.

NOTE... these are oscillating solutions... even though they are expressed as sums!