

Classification and Canonical Forms.

Goal - Classify Second order PDEs.

$$\text{General Form } Au_{xx} + Bu_{xy} + Cu_{yy} + Du_x + Eu_y + Fu = G$$

$$\left. \begin{array}{l} \text{Hyperbolic: } B^2 - 4AC > 0 \\ \text{Parabolic: } B^2 - 4AC = 0 \\ \text{Elliptic: } B^2 - 4AC < 0 \end{array} \right\} \begin{array}{l} \text{solution methods and} \\ \text{possible canonical forms} \\ \text{depend on Classification.} \end{array}$$

Ex $u_t = u_{xx} \longrightarrow u_{xx} - u_t = 0$ here $y = t$

$$A=1 \quad E=-1 \quad \text{all others zero} \\ 0^2 - 4(1)(0) = 0 \quad \text{PARABOLIC!}$$

Ex $u_{tt} = u_{xx} \longrightarrow u_{xx} - u_{tt} = 0$

$$A=1 \quad C=-1 \quad \text{all others zero} \\ 0^2 - 4(1)(-1) = 4 > 0 \quad \text{HYPERBOLIC!}$$

Ex $u_{xx} + u_{yy} = 0$

$$A=1 \quad B=0 \quad C=1 \quad \text{all others zero} \\ 0^2 - 4(1)(1) = -4 < 0 \quad \text{ELLIPTIC!}$$

Ex $xu_{xx} + u_{yy} = \sin(x)$ variable coefficients

$$A=x \quad C=1 \quad \text{all others zero} \\ 0^2 - 4(x)(1) = -4x \quad \begin{array}{l} x < 0 \text{ Hyperbolic} \\ x = 0 \text{ Parabolic} \\ x > 0 \text{ Elliptic} \end{array}$$

Each Equation Type has a specific CANONICAL FORM.

$$\text{HYPERBOLIC} \longrightarrow \begin{array}{l} u_{\xi\xi} - u_{\eta\eta} = \psi(\xi, \eta, u, u_\xi, u_\eta) \\ u_{\xi\eta} = \phi(\xi, \eta, u, u_\xi, u_\eta) \end{array}$$

$$\text{PARABOLIC} \longrightarrow u_{\eta\eta} = \phi(\xi, \eta, u, u_\xi, u_\eta)$$

$$\text{ELLIPTIC} \longrightarrow u_{\xi\xi} + u_{\eta\eta} = \phi(\xi, \eta, u, u_\xi, u_\eta)$$

Today - tricks for simplifying more complicated PDEs.

Nondimensionalization - often it's easier to deal with problems where the dimensions have been taken out, or normalized.

Transform seemingly different problems into the same form.

Ex. Diffusion Problem

$$\left. \begin{aligned} u_t &= \alpha^2 u_{xx} \\ u(0,t) &= T_1 \\ u(L,t) &= T_2 \\ u(x,0) &= \sin\left(\frac{\pi x}{L}\right) \end{aligned} \right\} \begin{aligned} &\text{Want to change the problem} \\ &1. \text{ No physical parameters} \\ &2. \text{ BCs and ICs are simpler} \end{aligned}$$

Do this through a change of variables:

$$u(x,t) \longrightarrow V(\xi, \tau).$$

First define a new Temperature.

$$V(x,t) = \frac{u(x,t) - T_1}{T_2 - T_1} \quad \text{The new "temperature" has no units.}$$

Use reference temperatures from the boundaries.

Transform Each Piece.

$$u(x,t) = (T_2 - T_1)V(x,t) + T_1$$

$$u_t = (T_2 - T_1)V_t \quad u_{xx} = (T_2 - T_1)V_{xx}$$

$$u(0,t) = (T_2 - T_1)V(0,t) + T_1 = T_1 \quad V(0,t) = 0$$

$$u(L,t) = (T_2 - T_1)V(L,t) + T_1 = T_2 \quad V(L,t) = 1$$

$$u(x,0) = (T_2 - T_1)V(x,0) + T_1 = \sin\left(\frac{\pi x}{L}\right) \quad V(x,0) = \frac{\sin\left(\frac{\pi x}{L}\right) - T_1}{T_2 - T_1}$$

NEW PDE

$$\left. \begin{aligned} U_t &= \alpha^2 U_{xx} & 0 < x < L \\ U(0, t) &= 0 \\ U(1, t) &= 1 \\ U(x, 0) &= \frac{\sin(\pi x/L) - T_1}{T_2 - T_1} \end{aligned} \right\}$$

NOTE... nondims are not unique. We could have picked $U(x, t) = \frac{u(x, t)}{T_2}$

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Transform space - choose a "length" variable with no units.

let $\xi = x/L$ * usually wrt the length of the domain of the diffusion length.

$$U_x = U_\xi \frac{d\xi}{dx} = \frac{U_\xi}{L} \quad \text{and} \quad U_{xx} = U_{\xi\xi} \left(\frac{d\xi}{dx}\right)^2 = \frac{U_{\xi\xi}}{L^2}$$

NEW PDE

$$\left. \begin{aligned} U_t &= \left(\frac{\alpha^2}{L^2}\right) U_{\xi\xi} & 0 < \xi < 1 \\ U(0, t) &= 0 \\ U(1, t) &= 1 \\ U(\xi, 0) &= \frac{\sin(\pi\xi) - T_1}{T_2 - T_1} \end{aligned} \right\}$$

the domain is now between zero and one.

Transform time - choose a "time" variable that simplifies everything else.

let $\tau = ct$ plug in and solve for c.

$$U_t = U_\tau \frac{d\tau}{dt} = c U_\tau$$

$$\text{so } c U_\tau = \frac{\alpha^2}{L^2} U_{\xi\xi} \quad \text{and} \quad U_\tau = \frac{1}{c} \frac{\alpha^2}{L^2} U_{\xi\xi}$$

$$\text{so } c = \alpha^2/L^2 \quad \text{making } \tau = \frac{\alpha^2}{L^2} t$$

NEW PDE

$$\left. \begin{aligned} U_\tau &= U_{\xi\xi} & 0 < \xi < 1 \\ U(0, \tau) &= 0 \\ U(1, \tau) &= 1 \\ U(\xi, 0) &= \phi(\xi) \end{aligned} \right\}$$

much simpler and more compact... probably a problem we have solved before!

- After finding the answer we transform back.

You Try...

$$u_{tt} = \alpha^2 u_{xx}$$

$$u(0, t) = 0$$

$$u(L, t) = 0$$

$$u(x, 0) = \sin(\pi x/L) + \frac{1}{2} \sin(3\pi x/L)$$

$$u_t(x, 0) = 0$$

Note transform
only

$$x \rightarrow \xi$$

$$t \rightarrow \tau$$

because the
BC's are already
simple.

let $\xi = x/L$

Then $\frac{u_{\xi\xi}}{L^2} = u_{xx}$

NEW PDE

$$u_{tt} = \alpha^2/L^2 u_{\xi\xi}$$

$$u(0, t) = 0$$

$$u(1, t) = 0$$

$$u(\xi, 0) = \sin(\pi \xi) + \frac{1}{2} \sin(3\pi \xi)$$

$$u_t(\xi, 0) = 0$$

Then solve by recognizing
the problem from
Lesson 20 prob 1

then let $\tau = ct$ and $u_{tt} = c^2 u_{\tau\tau}$

$$c^2 u_{tt} = \alpha^2/L^2 u_{\xi\xi} \rightarrow u_{\tau\tau} = \frac{1}{c^2} \frac{\alpha^2}{L^2} u_{\xi\xi}$$

This means that $c = \alpha/L$
 $\tau = \alpha/L t$

NEW PDE.

$$u_{\tau\tau} = u_{\xi\xi}$$

$$u(0, \tau) = 0$$

$$u(1, \tau) = 0$$

$$u(\xi, 0) = \sin(\pi \xi) + \frac{1}{2} \sin(3\pi \xi)$$

$$u_t(\xi, 0) = 0$$

we solved this... $u(\xi, \tau) = \cos(\pi \tau) \sin(\pi \xi) + \frac{1}{2} \cos(3\pi \tau) \sin(3\pi \xi)$

giving the full soln:

$$u(x, t) = \cos\left(\frac{\pi \alpha t}{L}\right) \sin\left(\frac{\pi x}{L}\right) + \frac{1}{2} \cos\left(\frac{3\pi \alpha t}{L}\right) \sin\left(\frac{3\pi x}{L}\right)$$

Recall in The D'Alembert Soln.

$$u_{tt} = c^2 u_{xx} \text{ we introduced } \begin{matrix} \xi = x+ct \\ \eta = x-ct \end{matrix} \longrightarrow u_{\xi\eta} = 0$$

we were switching to canonical form!!

* Doing this either greatly simplifies the PDE
or sets us up for straight forward numerical
analysis !!

Canonical Form - HYPERBOLIC ... want $u_{\xi\eta} = \phi(\xi, \eta, u, u_\xi, u_\eta)$

1. Define a coordinate transformation

$$\xi = \xi(x, t) \quad \eta = \eta(x, t) \longrightarrow \text{goal of the method is to find these !!}$$

Consider the general Form

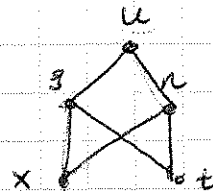
$$A u_{xx} + B u_{xt} + C u_{tt} + D u_x + E u_t + F u = G$$

and sub in our new coordinates:

$$u = u$$

$$u_x = u_\xi \xi_x + u_\eta \eta_x$$

$$u_t = u_\xi \xi_t + u_\eta \eta_t$$



now the second derivatives.

$$u_{xx} = \frac{\partial}{\partial x} [u_x] = \frac{\partial}{\partial x} [u_3 z_x + u_n n_x]$$

$$= u_{33} (z_x)^2 + u_{33x} z_x + u_{3n} z_x n_x + u_{n3} n_x z_x + u_{nn} (n_x)^2 + u_{nnx} n_x$$

$$= u_{33} (z_x)^2 + 2u_{3n} z_x n_x + u_{nn} (n_x)^2 + u_{33x} z_x + u_{nnx} n_x.$$

similarly... on page 177 !!

$$u_{xt} = u_{33} z_x z_t + u_{3n} (z_x n_t + z_t n_x) + u_{nn} n_x n_t + u_{33t} z_x + u_{nnx} n_x$$

$$u_{tt} = u_{33} z_t^2 + 2u_{3n} z_t n_t + u_{nn} n_t^2 + u_{33t} z_t + u_{nnx} n_t$$

We substitute these into the general form to get:

$$\bar{A} u_{33} + \bar{B} u_{3n} + \bar{C} u_{nn} + \bar{D} u_3 + \bar{E} u_n + \bar{F} u = \bar{G}$$

where

$$\bar{A} = A z_x^2 + B z_x z_t + C z_t^2$$

$$\bar{B} = 2A z_x n_x + B (z_x n_t + z_t n_x) + 2C z_t n_t$$

$$\bar{C} = A n_x^2 + B n_x n_t + C n_t^2$$

$$\bar{D} = A z_{xx} + B z_{xt} + C z_{tt} + D z_x + E z_t$$

:

again in The Book.

But NOW WE SET $\bar{A} = 0$ and $\bar{C} = 0$ to force the canonical Form.

This leads to

$$A \left[\frac{z_x}{z_t} \right]^2 + B \left[\frac{z_x}{z_t} \right] + C = 0$$

$$A \left[\frac{n_x}{n_t} \right]^2 + B \left[\frac{n_x}{n_t} \right] + C = 0$$

quadratic
equations for
 z_x/z_t and n_x/n_t

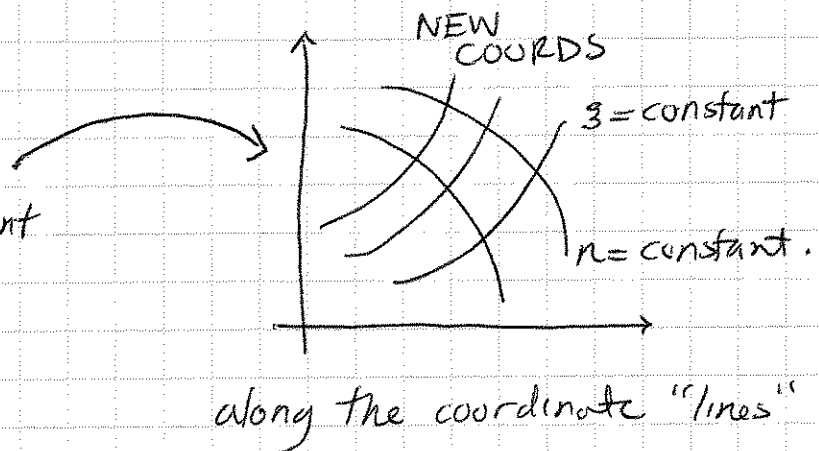
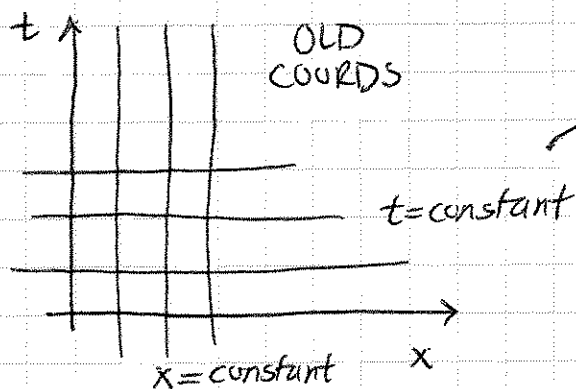
where $\left[\frac{z_x}{z_t} \right] = \frac{-B + \sqrt{B^2 - 4AC}}{2A}$

$\left[\frac{n_x}{n_t} \right] = \frac{-B - \sqrt{B^2 - 4AC}}{2A}$

note we just need one soln for each that will work and since we want different coordinates we choose one positive and one negative.

This is also where $B^2 - 4AC$ definitions come from. the idea of the characteristics.

NOW HOW DO WE RELATE THIS TO OUR ORIGINAL x, t COORDINATES.



EX $U_{xx} - 4U_{tt} + U_x = 0$

$\left[\frac{z_x}{z_t} \right] = 2$ and $\left[\frac{n_x}{n_t} \right] = 2$

$\frac{dz}{dx} = -2$ $\frac{dn}{dx} = 2$

integrate

$z = -2x + C_1$ $n = 2x + C_2$

OR
 $z + 2x = C_1$ $n - 2x = C_2$

$z = t + 2x = C_1$ $n = t - 2x = C_2$

$z = \text{constant}$ $n = \text{constant}$
 $dz = z_x dx + z_t dt = 0$ $dn = n_x dx + n_t dt = 0$

$\frac{dt}{dx} = -\left[\frac{z_x}{z_t} \right]$ OR $\frac{dt}{dx} = -\left[\frac{n_x}{n_t} \right]$

we can use these to solve for z and n in terms of x and t .

So we can transform to canonical form using

$z = t + 2x$ and $n = t - 2x$!!